

Recent shell-model studies of nuclear structure and applications to radiative neutron capture

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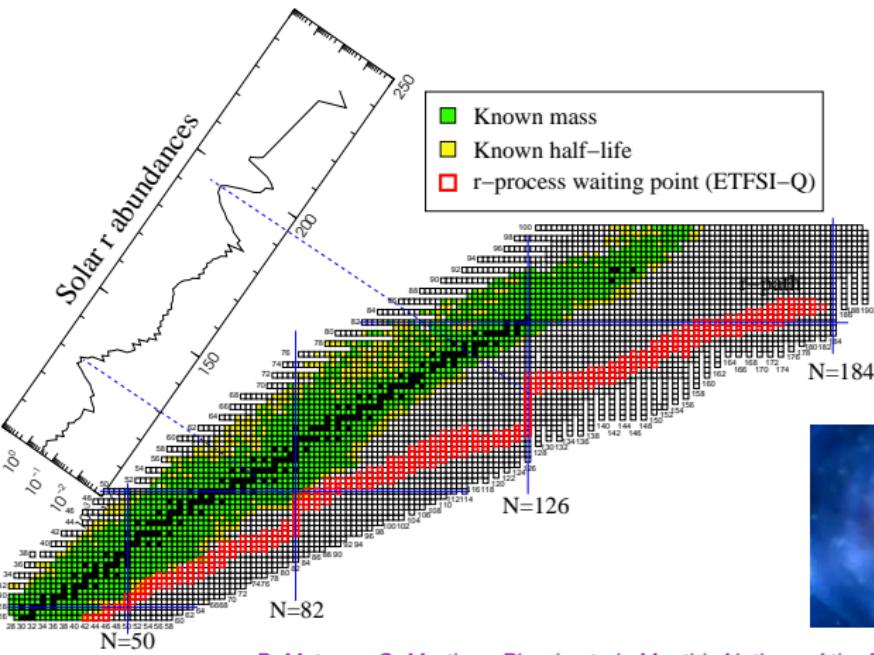


IN2P3 theory project: "Photoabsorption Strength Functions from CI method"

NUSPIN, 25.06.2019

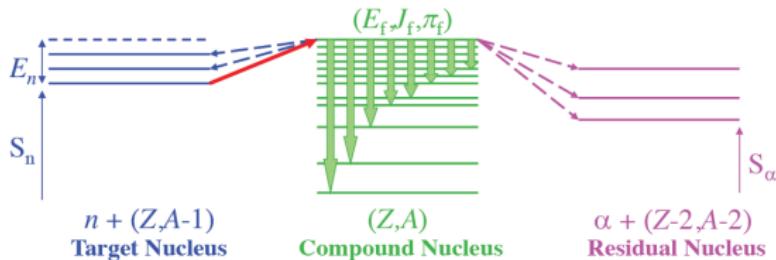
R-process nucleosynthesis

- Nuclear models are needed to provide input for r-process simulations: masses, *level densities (ρ)*, β half-lives, γ -Strength Functions (f_{XL}), fission barriers...



B. Metzger, G. Martinez-Pinedo et al., Monthly Notices of the Royal Astronomical Society. 406, 2650 (2010)

Radiative neutron capture: resonant capture



$$\sigma_{(n,\gamma)}^{\mu\nu}(E_i, n) = \frac{\pi\hbar^2}{2M_{i,n}E_{i,n}} \frac{1}{(2J_i^\mu+1)(2J_n+1)} \sum_{J,\pi} (2J+1) \frac{T_n^\mu T_\gamma^\nu}{T_n^\mu + T_\gamma^\nu}$$

for $E_n \sim \text{keV}$ $T_n^\mu \gg T_\gamma^\nu \rightarrow \sigma^{\mu\nu} \sim T_\gamma^\nu$

$E_{i,n}, M_{i,n}$ - center-of-mass energy, reduced mass of the system

$J_n = 1/2$ -neutron spin

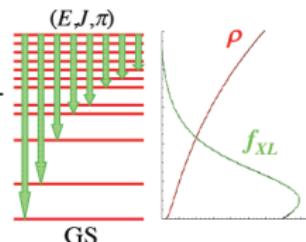
$T_n^\mu = T_n(E, J, \pi; E_i^\mu, J_i^\mu, \pi_i^\mu)$ $T_\gamma^\nu = T_\gamma(E, J, \pi; E_m^\nu, J_m^\nu, \pi_m^\nu)$ - transmission coefficients

For a given multipolarity

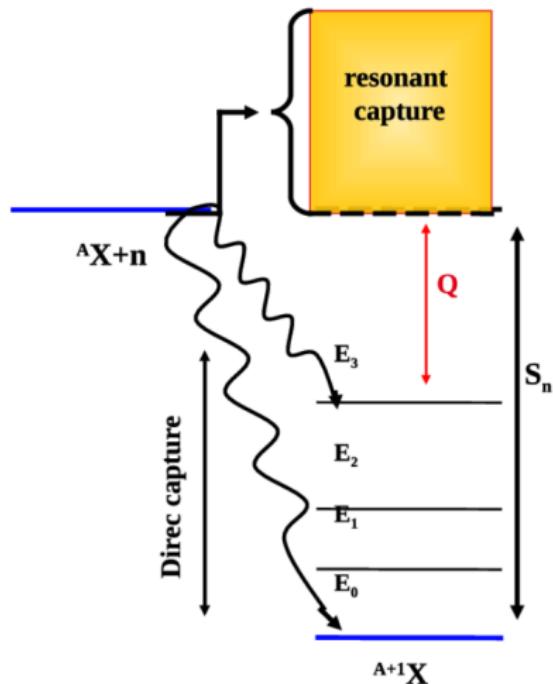
$$T_{XL}(E, J, \pi, E^\nu, J^\nu, \pi^\nu) = 2\pi E_\gamma^{2L+1} f_{XL}(E, E_\gamma)$$

Test, using SM, the key ingredients of Hauser-Feshbach calculations:

- description of γ emission spectra
- Brink-Axel hypothesis



Radiative neutron capture: direct capture



Xi. Yu and S. Goriely, Phys. Rev. C86 (2012) 045801

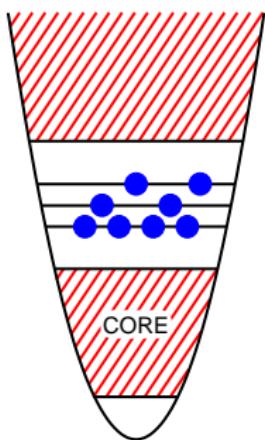
$$\sigma^{DC}(E) = \sum_{f=0}^x S_f \sigma_{dis}(E) + \langle S \rangle \int_{E_x}^{S_n} \sum_{J_f, \pi_f} \rho(E_f, J_f, \pi_f) \times \sigma_f^{cont} dE_f$$

If no experimental data available:

- use combinatorial model for the level density with $\langle S \rangle = 0.5$
 - The key ingredients: low-energy levels and spectroscopic factors
 - Validate theoretical approximations (HFB) in exotic nuclei using SM predictions

Shell model: generalities

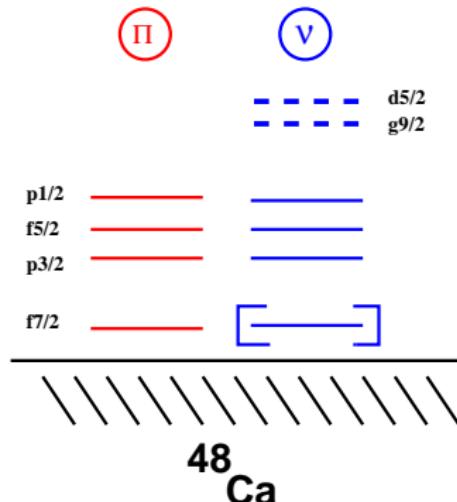
Shell model relies on the possibility of diagonalizing the Hamiltonian matrix and deriving (constraining empirically) a suitable effective interaction.



$$H_{\text{eff}} |\Psi_{\text{eff}}\rangle = E |\Psi_{\text{eff}}\rangle$$

- **Direct capture:** knowledge of the lowest-lying levels (energies and spectroscopic factors) → quality of the effective Hamiltonian
- **Resonant capture:** knowledge of statistical properties (energies and transitions in nuclear continuum) → possibility of computing of hundreds of nuclear levels

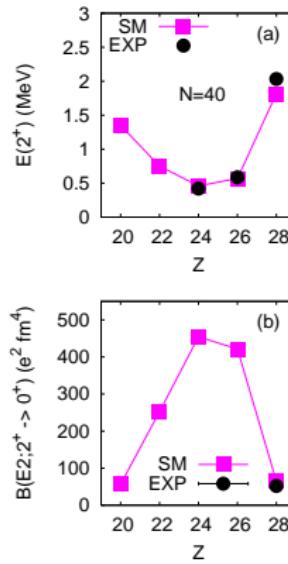
Lenzi-Nowacki-Poves-Sieja interaction



- based on realistic TBME
- new fit of the pf shell (KB3GR, E. Caurier)
- monopole corrections

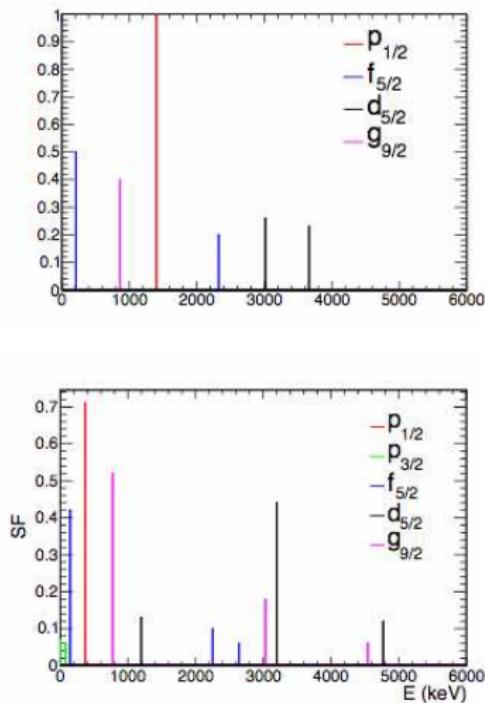
S. Lenzi et al, Phys. Rev. C82 (2010) 054301

Used successfully in over 50 papers (masses, spectroscopy, transitions, spectroscopic factors...)

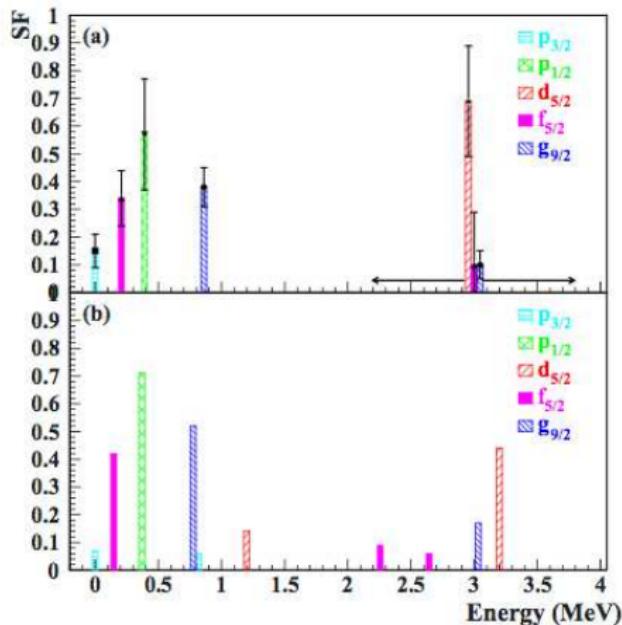


^{60}Fe spectroscopic factors

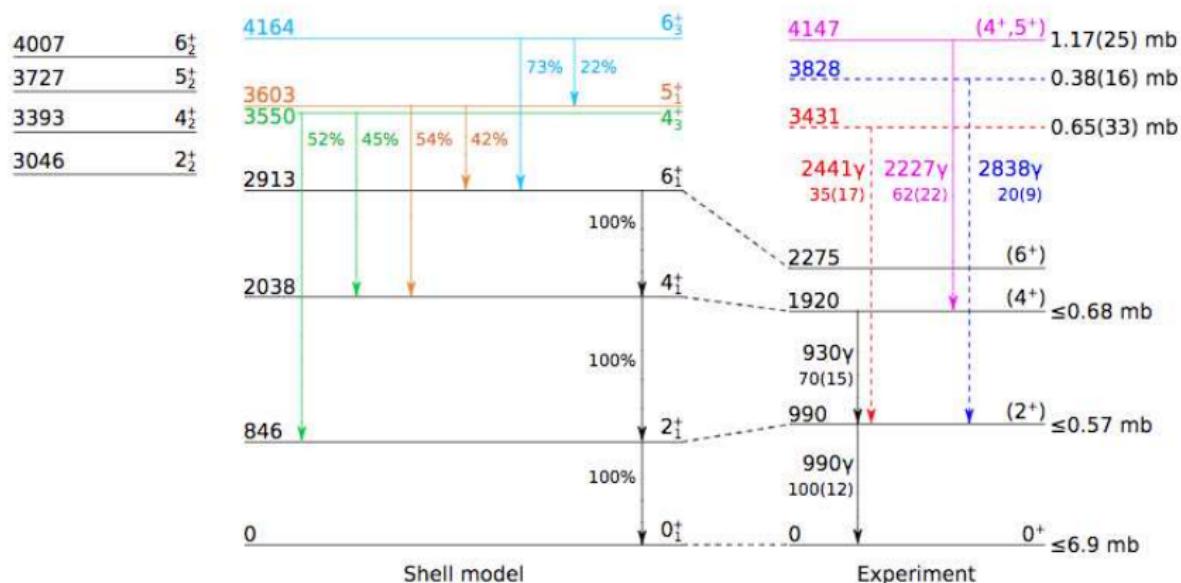
S. Giron, PhD Orsay, 2012



S. Giron et al., Phys. Rev. C95 (2017) 035806



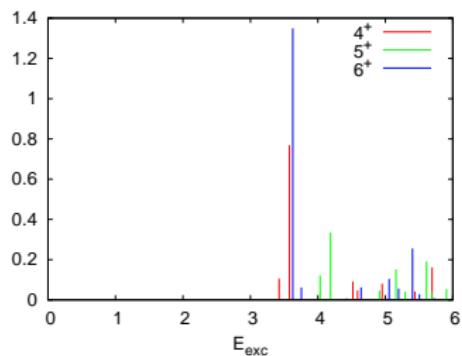
^{76}Ni from ($p, 2p$) reaction



Z. Elekes et al., Phys. Rev. C99 (2019) 014312

^{76}Ni from ($p, 2p$) reaction

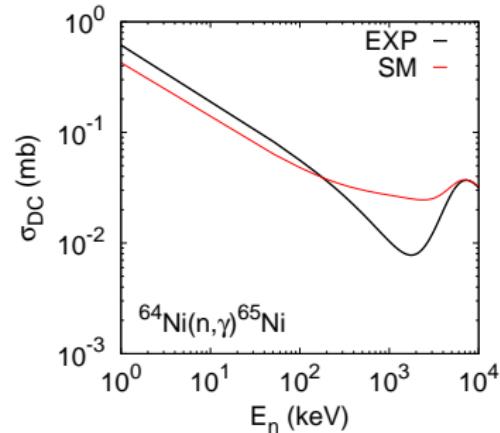
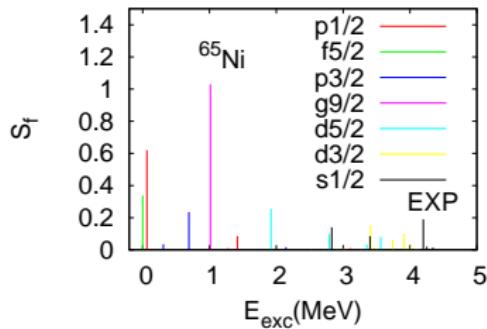
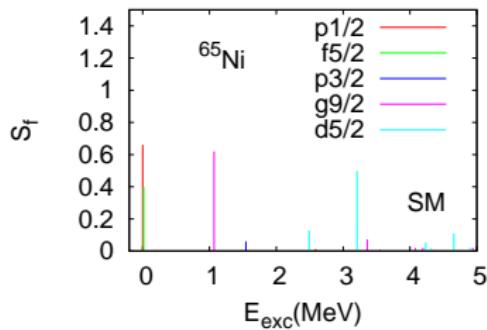
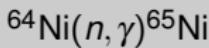
S_{7/2}



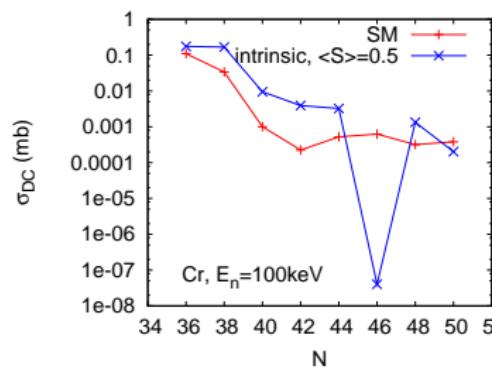
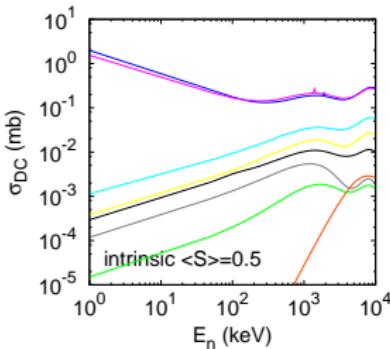
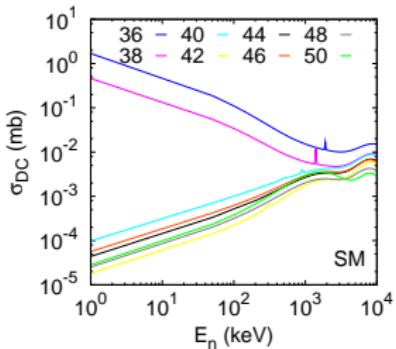
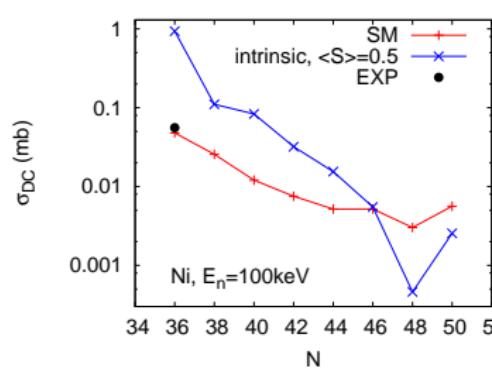
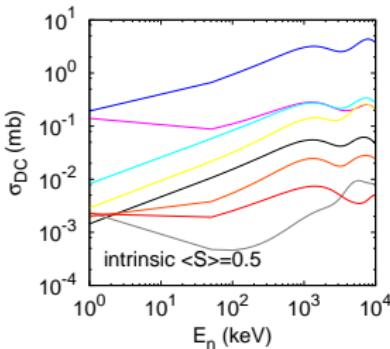
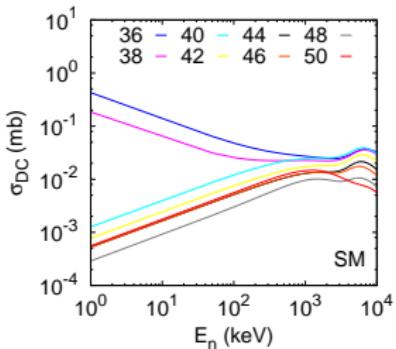
- Proton gap $Z = 28$ estimated to be 4.4MeV in ^{76}Ni and 5.2MeV in ^{78}Ni , compatible with other recent studies in the region.

E_{theor} (MeV)	J^π	Spectroscopic factor	σ_{theor} (mb)	E_{expt} (MeV)	σ_{expt} (mb)
0	0_1^+	0.60	$f_{5/2}$	0.74	0
0.846	2_1^+	<0.01	$f_{7/2}$	0.28	0.990
			$f_{5/2}$	0.23	
2.038	4_1^+	<0.01	$f_{7/2}$	<0.10	1.920
		<0.01	$f_{5/2}$		
2.913	6_1^+	<0.01	$f_{7/2}$	<0.10	2.275
					3.431
					3.828
					0.65(33)
3.550	4_3^+	0.77	$f_{7/2}$	1.08	4.147
		<0.01	$f_{5/2}$		1.17(25)
3.603	5_1^+	0.72	$f_{7/2}$	1.02	
		<0.01	$f_{5/2}$		
4.164	6_3^+	0.62	$f_{7/2}$	0.86	

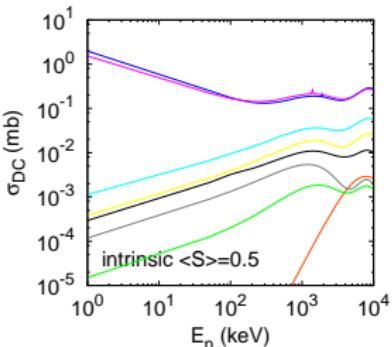
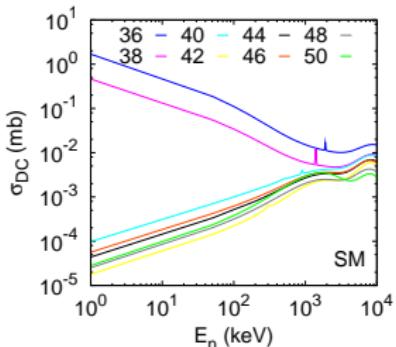
Direct capture rates using SM results



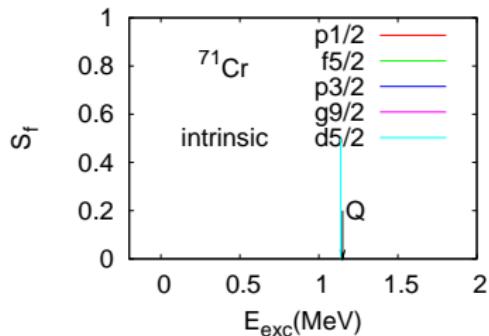
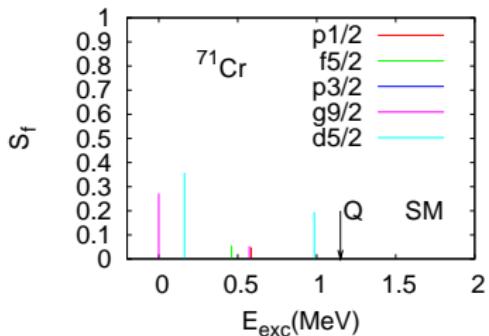
Direct capture rates using SM: Ni & Cr



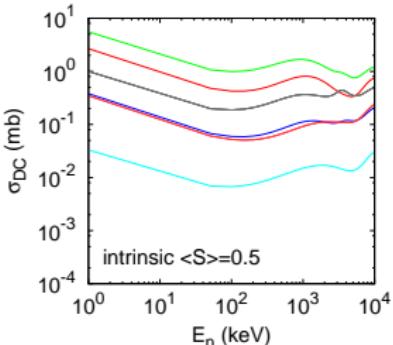
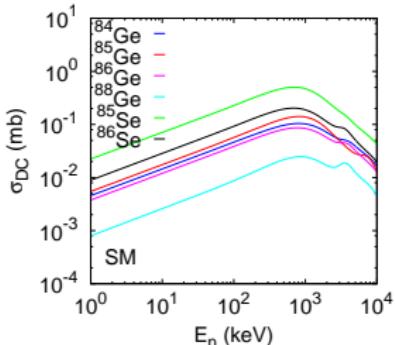
Direct capture rates using SM results: chromiums



Great sensitivity to the structure details in very neutron-rich nuclei.



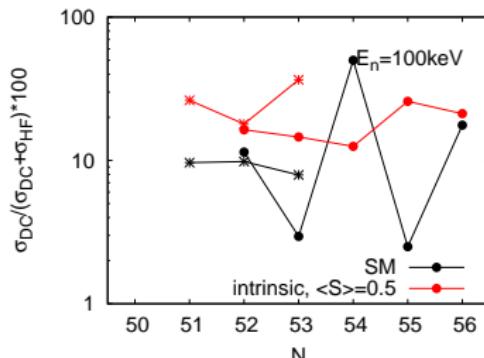
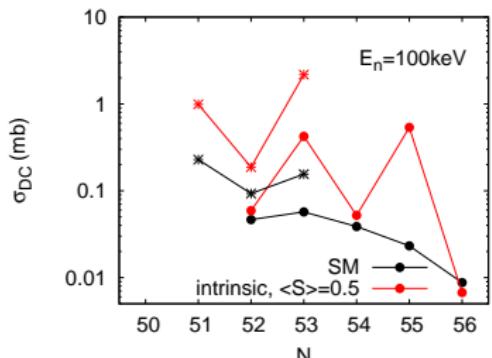
Direct capture using SM: triaxial nuclei above ^{78}Ni



Calculations with NI78-II interaction
Prediction of triaxiality above
 ^{78}Ni

*K. Sieja et al., Phys. Rev. C88 (2013)
034327*

*M. Lettmann et al., Phys. Rev. C96
(2017) 011301R*



■ Spectroscopy of odd- N nuclei past $N = 50$ of interest for nuclear models

Going beyond usual SM applications: γ -decay

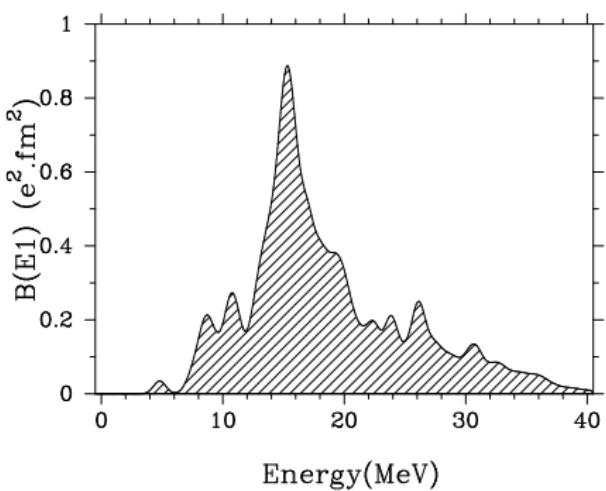
photoabsorption (PSF)

- use Lanczos Strength Function method with a large number of iterations

EXP: $\sum B(E1) = 0.49 \pm 0.16 \text{ e}^2\text{fm}^2$ (6-10MeV)

J. Gobel et al., Phys. Rev. Lett. 101 (2008) 212503

THEO: $\sum B(E1) = 0.485 \text{ e}^2\text{fm}^2$ (0-10MeV)

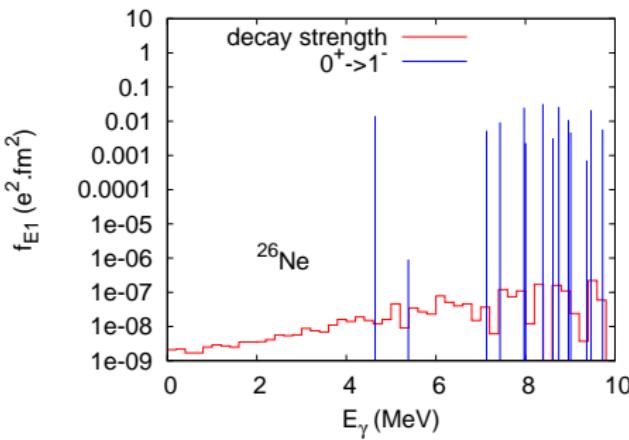


γ -decay (RSF)

- calculate desired number of low lying states using standard SM diagonalization techniques
- obtain the averages and radiative strength functions from relations:

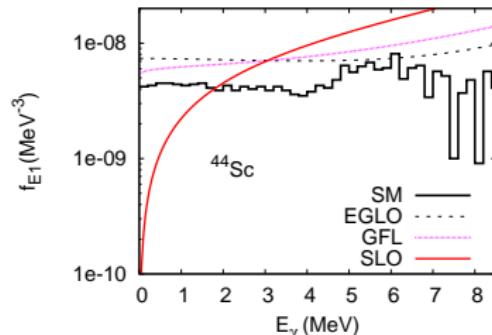
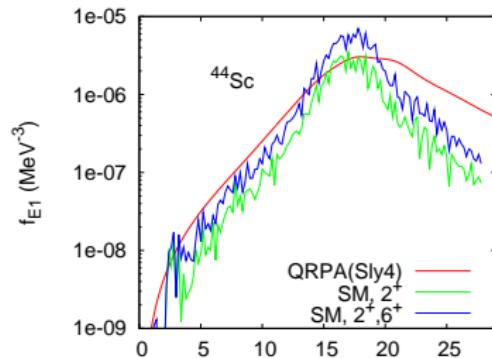
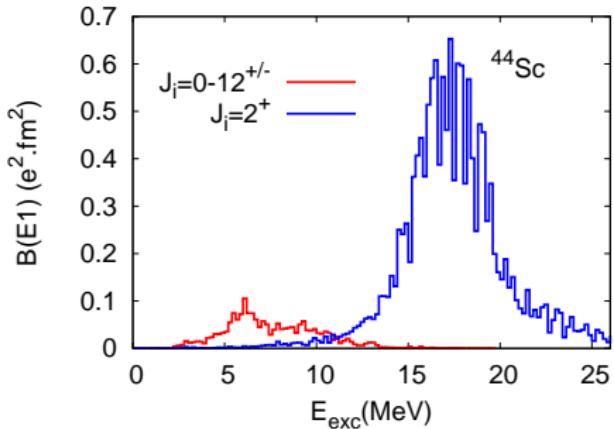
$$f_{M1/E1}(E_\gamma) = 16\pi/9(\hbar c)^3 S_{M1/E1}(E_\gamma)$$

$$S_{M1/E1} = \langle B(M1/E1) \rangle \rho_i(E_i)$$



$E1$ strength: comparison to other models

K. Sieja, PRL119 (2017) 052502



- Description of the resonance and radiative decay in the same theoretical framework.
- Reasonable agreement between QRPA and SM PSF up to 15 MeV.
- Microscopic SM strength has a non-zero limit for $E_\gamma = 0$. Consistent with the EGLO model.

Dipole strength in ^{44}Sc : theory vs exp

PRL 119, 052502 (2017)

PHYSICAL REVIEW LETTERS

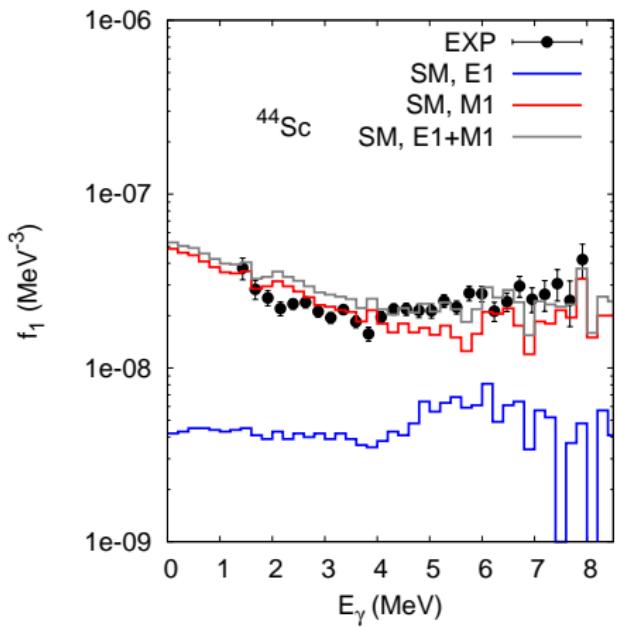
week ending
4 AUGUST 2017

Electric and Magnetic Dipole Strength at Low Energy

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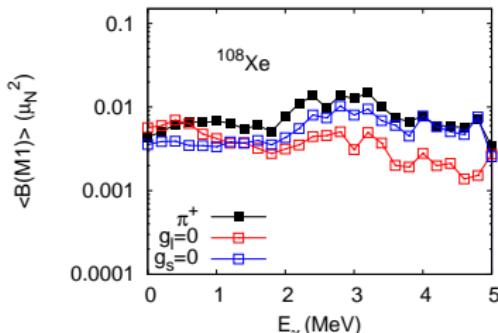
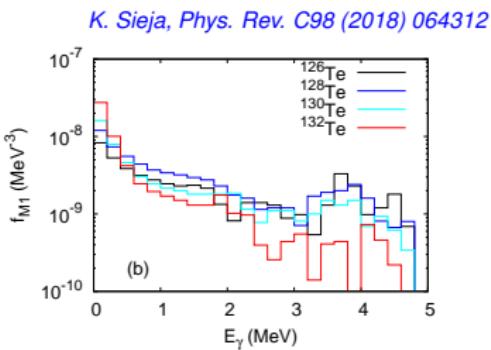
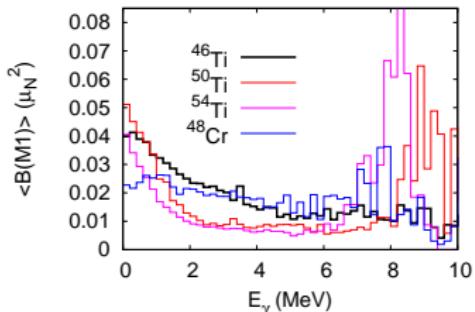
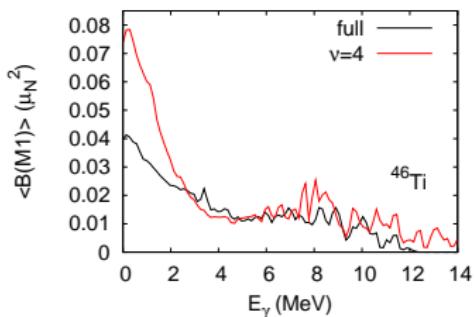
- all states below $S_n \sim 10\text{MeV}$
- 86642 $M1$ matrix elements
- 65670 $E1$ matrix elements

- Overall good agreement with data:
- the upbend is due to $M1$ transitions
- the $E1$ pattern is flat with the non-zero $E_\gamma \rightarrow 0$ limit

M1 upbend: general trends

K. Sieja, Eur. Phys. J. Web of Conf. 146 (2017) 05004 ND2016

S. Karampagia et al., Phys. Rev. C95 (2017) 024322

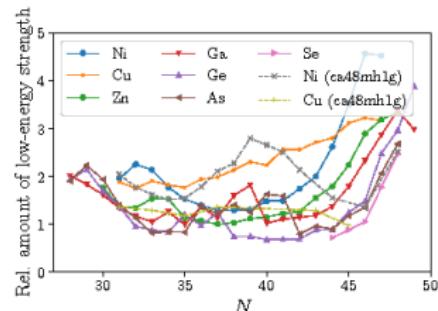


- The strength at $E_\gamma = 0$ peaks around shell closures and is flat in deformed nuclei

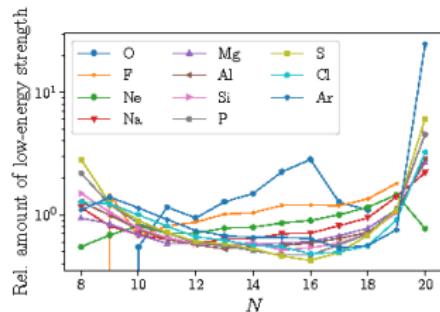
Shell structure in the nuclear quasi-continuum

Ratio of $B(M1)$ strength: $\frac{\langle B(M1) \rangle(0-2\text{MeV})}{\langle B(M1) \rangle(2-6\text{MeV})}$

mid-mass nuclei:

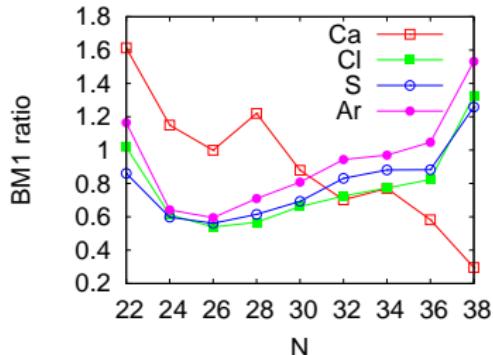


sd-shell nuclei:



J. Midtbo et al., Phys. Rev. C98 (2018) 064321

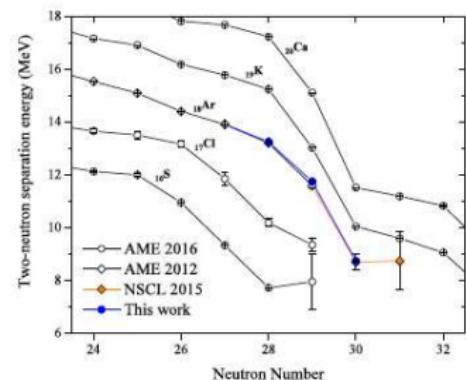
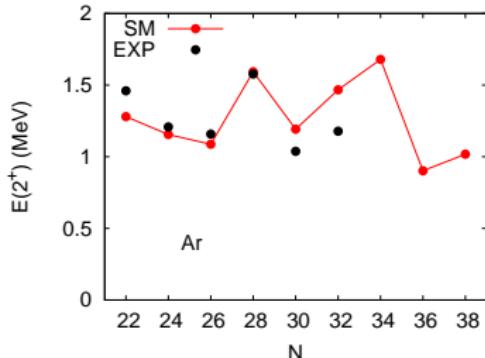
sd – pf nuclei:



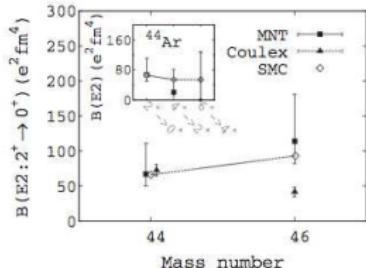
- The ratio peaks towards the edges of the model spaces at $N = 8, 20, 40, 50$.
- Some extra shell effects are present in the Ni chain.
- In sd-pf nuclei the ratio peaks at $N = 28$ for Ca only.

The puzzling case of ^{46}Ar

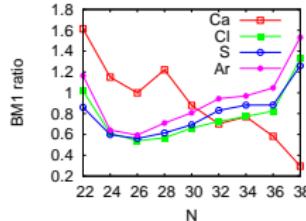
D. Mengoni et al., Phys. Rev. C82 (2010) 024308



M. Maugeot, PhD thesis, Orsay 2018

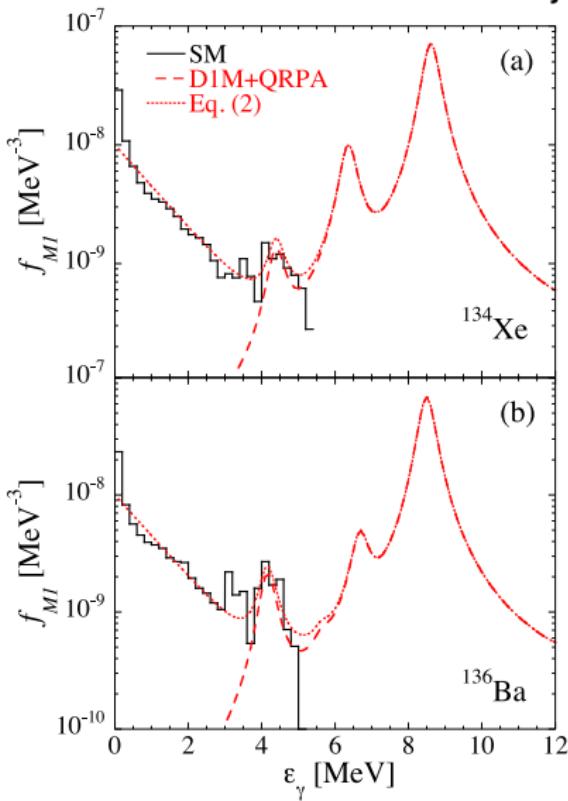


- Signs of shell closure in ^{46}Ar from 2^+ energies and two-neutron separation energies
- Confusing evidence from $B(E2)$ values
- No shell closure from the $B(M1)$ decay strength!



Application of the low-energy limit to the QRPA results

Collaboration with S. Goriely (ULB), S. Hilaire and S. Péru (CEA-DAM)



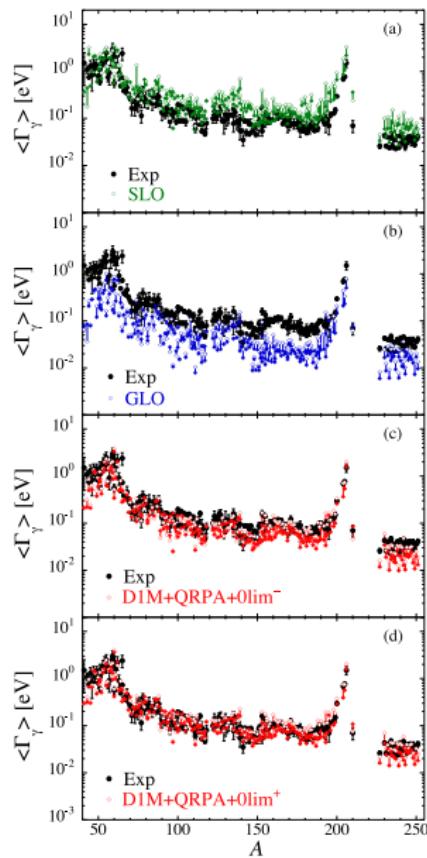
S. Goriely, S. Hilaire, S. Péru and K. Sieja, PRC98 (2018) 014327

To describe radiative decay, phenomenological low-energy corrections fitted to reproduce SM trends and data are added to microscopic QRPA-Gogny $M1$ and $E1$ PSF:

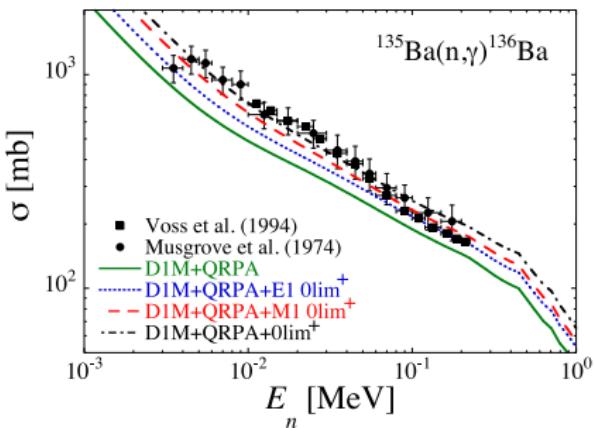
$$\begin{aligned} f_{E1}(\varepsilon_\gamma) &= f_{E1}^{QRPA}(\varepsilon_\gamma) + f_0 U/[1 + e^{(\varepsilon_\gamma - \varepsilon_0)}] \\ f_{M1}(\varepsilon_\gamma) &= f_{M1}^{QRPA}(\varepsilon_\gamma) + C e^{-\eta \varepsilon_\gamma} \end{aligned} \quad (2)$$

- upper limit ($0\lim^+$)
 $f_0 = 5 \cdot 10^{-10} \text{ MeV}^{-4}$, $\varepsilon_0 = 5 \text{ MeV}$,
 $C = 3 \cdot 10^{-8} \text{ MeV}^{-3}$, $\eta = 0.8 \text{ MeV}^{-1}$
- lower limit ($0\lim^-$)
 $f_0 = 10^{-10} \text{ MeV}^{-4}$, $\varepsilon_0 = 3 \text{ MeV}$,
 $C = 10^{-8} \text{ MeV}^{-3}$, $\eta = 0.8 \text{ MeV}^{-1}$

Impact on radiative neutron capture

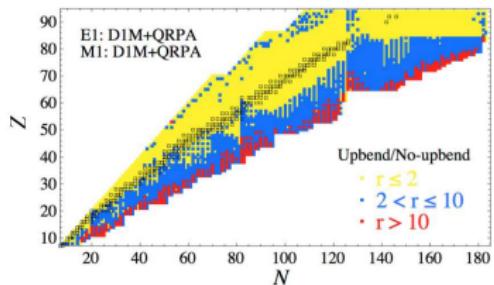
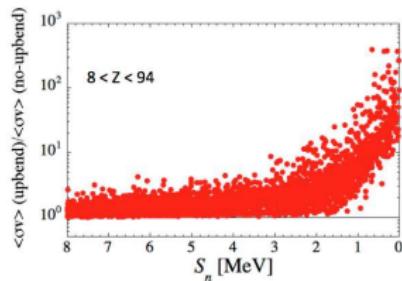
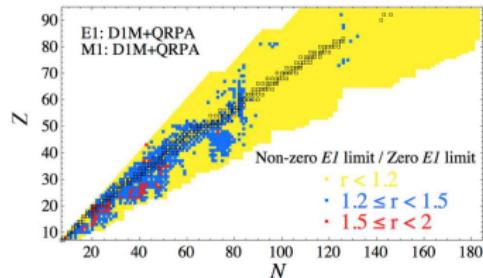
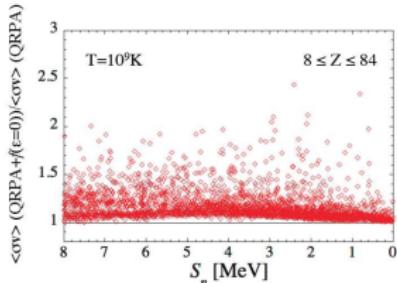


	rms/datum	
	$\langle \Gamma_\gamma \rangle$	$\langle \sigma \rangle$
0lim ⁻ (Comb)	0.88	1.07
0lim ⁻ (CT)	0.74	0.95
0lim ⁺ (Comb)	1.02	1.30
0lim ⁺ (CT)	0.90	1.15
GLO(Comb)	0.48	0.61
GLO(CT)	0.38	0.53



Impact on the radiative capture

MACS ratio at $T = 10^9$ K



- Non-zero limit of the $E1$ strength from SM has small impact on neutron capture: 20 – 50%
- $M1$ upbend can alternate the cross-section by a factor >10 in exotic nuclei

S. Goriely, S. Hilaire, S. Péru and K. Sieja, PRC98 (2018) 014327

Summary

- SM can provide reliable spectroscopic factors and help testing usual theoretical assumptions in cases no experimental data is known → work in progress.
- Spectroscopy of neutron-rich nuclei around ^{78}Ni is still of interest for nuclear models.
- $E1/M1$ RSF and PSF can be microscopically obtained within the SM.
- Shell effects survive at higher excitation energies and are visible in $M1$ dipole strength functions.
- $M1$ upbend has a significant impact on neutron capture cross sections in exotic nuclei: $\times 10$.